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THE AIR FORCE RESEARCH LABORATORY
LEAD | DISCOVER | DEVELOP | DELIVER

Kinetic Energy-Preserving Discretization Schemes for High Reynolds Number Propulsive Applications

Ayaboe Edoh, Ann Karagozian, Charles Merkle and Venke Sankaran

66th Annual APS Meeting
Fluid Dynamics Division

Pittsburgh, PA

Nov 24-26, 2013

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Objectives

Investigate dispersion and dissipation of numerical schemes with ultimate application to high-Re reacting LES

↓

Schemes

Standard Collocated Grid Schemes
Standard Staggered Grid Schemes
Kinetic Energy Preserving Schemes

↓

Analysis

Von Neumann Stability Analysis
1D Periodic Test Problem

Scope

$$\frac{\partial u}{\partial t} + a\frac{\partial u}{\partial x} = 0$$ Wave Eqn

$$\frac{\partial Q}{\partial t} + \frac{\partial E}{\partial x} = 0$$ Euler Eqns

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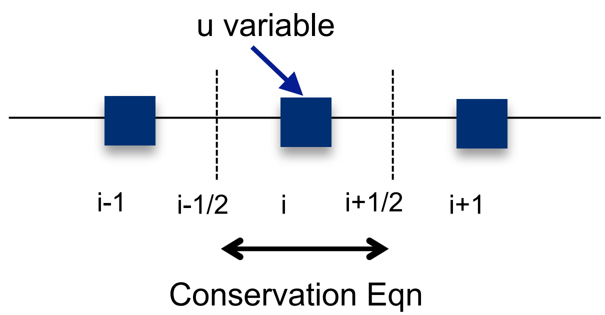
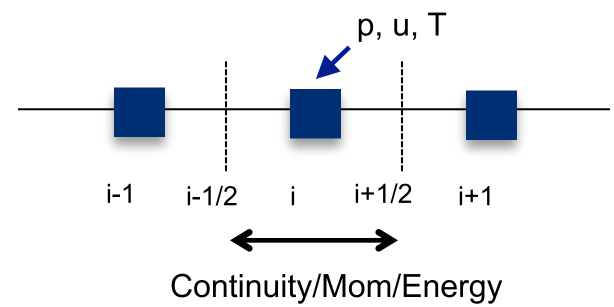
Formulation

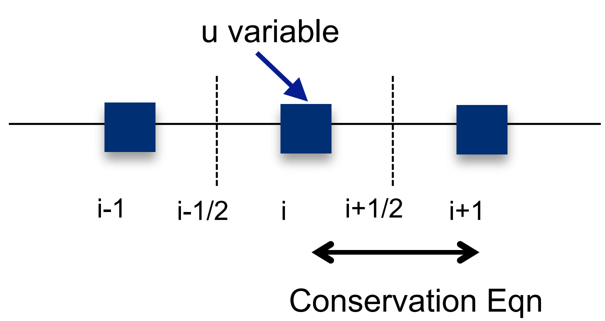
Wave Eqn / Euler Eqns

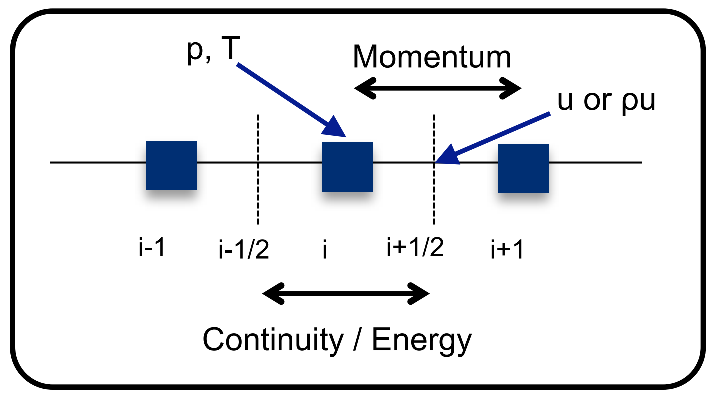
Collocated





Staggered





Variables also staggered in time for fully ke preserving schemes

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Von Neumann Analysis

Eigenvalues of the amplification matrix specify growth factor and phase errors.

$$Q^{n+1} = GQ^n$$

Staggered Grid Scheme

$$\Gamma_{ce}\left(\frac{\partial Q_{pT}}{\partial t}+\frac{\partial Q_u}{\partial t}\right)_i+\Gamma_m\left(\frac{\partial Q_{pT}}{\partial t}+\frac{\partial Q_u}{\partial t}\right)_{i+1/2}+A_{ce}\left(\frac{\partial Q_{pT}}{\partial x}+\frac{\partial Q_u}{\partial x}\right)_i+A_m\left(\frac{\partial Q_{pT}}{\partial x}+\frac{\partial Q_u}{\partial x}\right)_{i+1/2}=0$$

$$Q_{pT} = \begin{pmatrix} p \\ 0 \\ T \end{pmatrix}$$

Continuity/Energy

Momentum

$$Q_u = \begin{pmatrix} 0 \\ u \\ 0 \end{pmatrix}$$

Growth Factor

$$||g_i||$$

Phase Error

$$\frac{\phi}{\phi_{exact}} = \frac{-\tan^{-1}\{Imag(g_i)/Re(g_i)\}}{CFL \times \beta}$$

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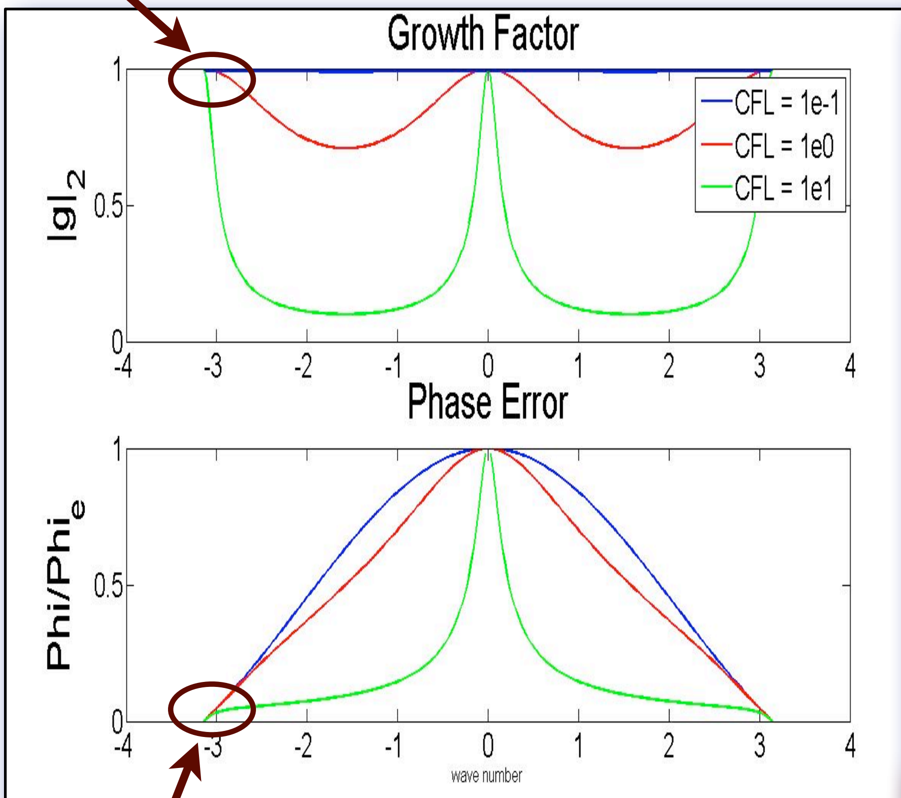
Wave Equation

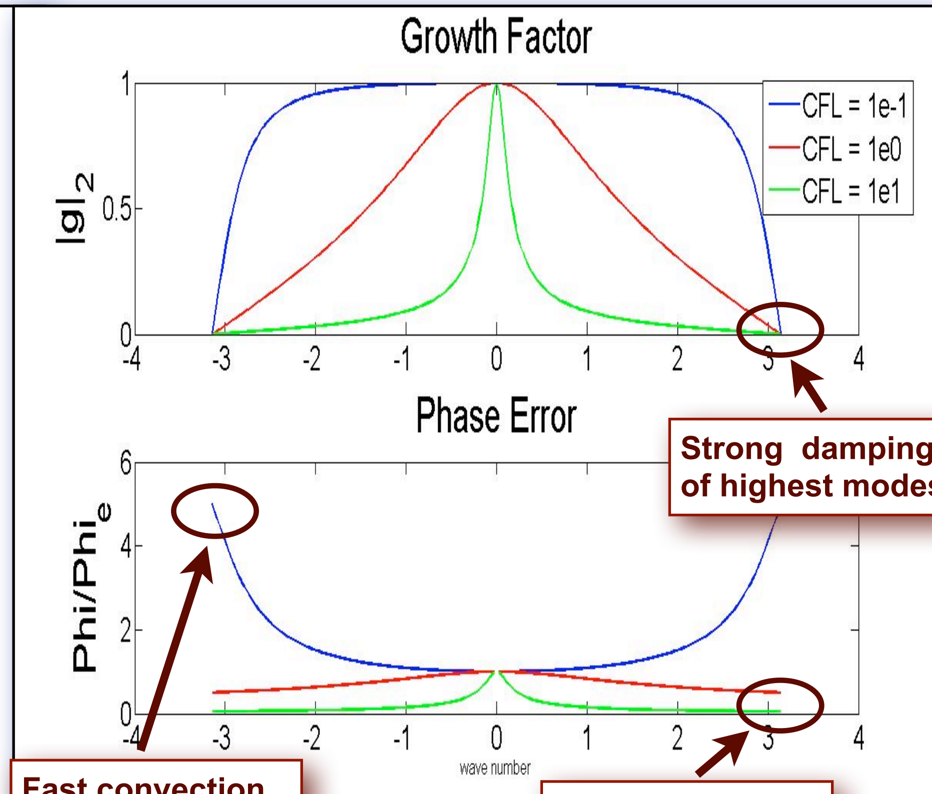
Euler Implicit Scheme

No damping of highest modes

Collocated

Staggered





Strong damping of highest modes

No convection of highest modes

Fast convection of highest modes for low CFL

Slow convection of highest modes for high CFL

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Euler Equations

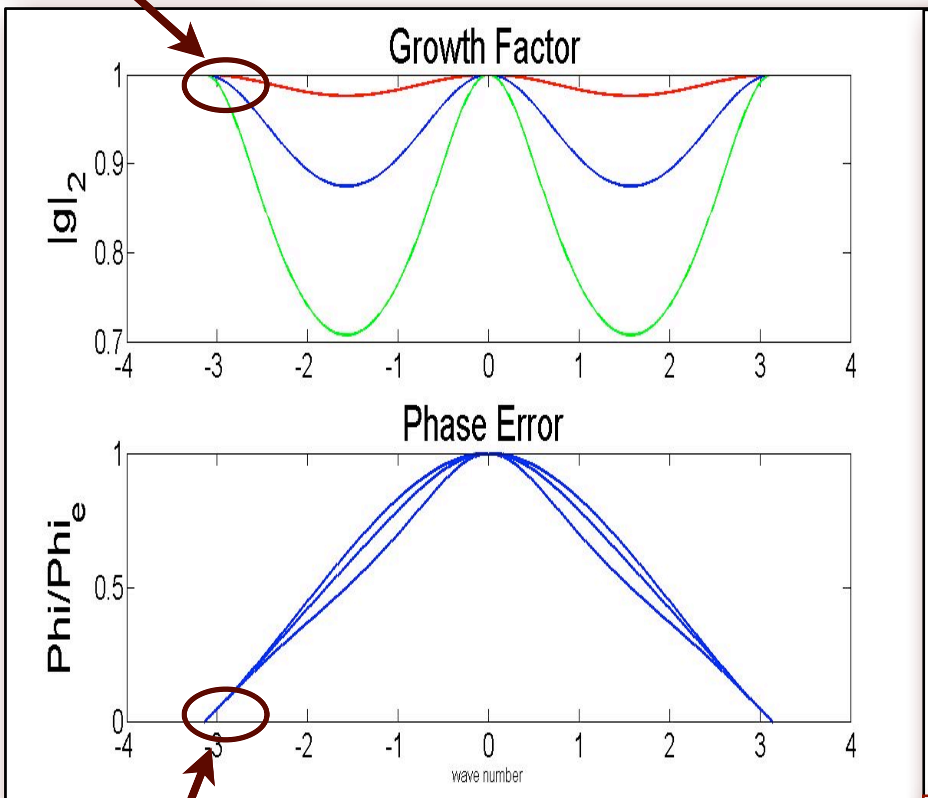
Euler Implicit Scheme

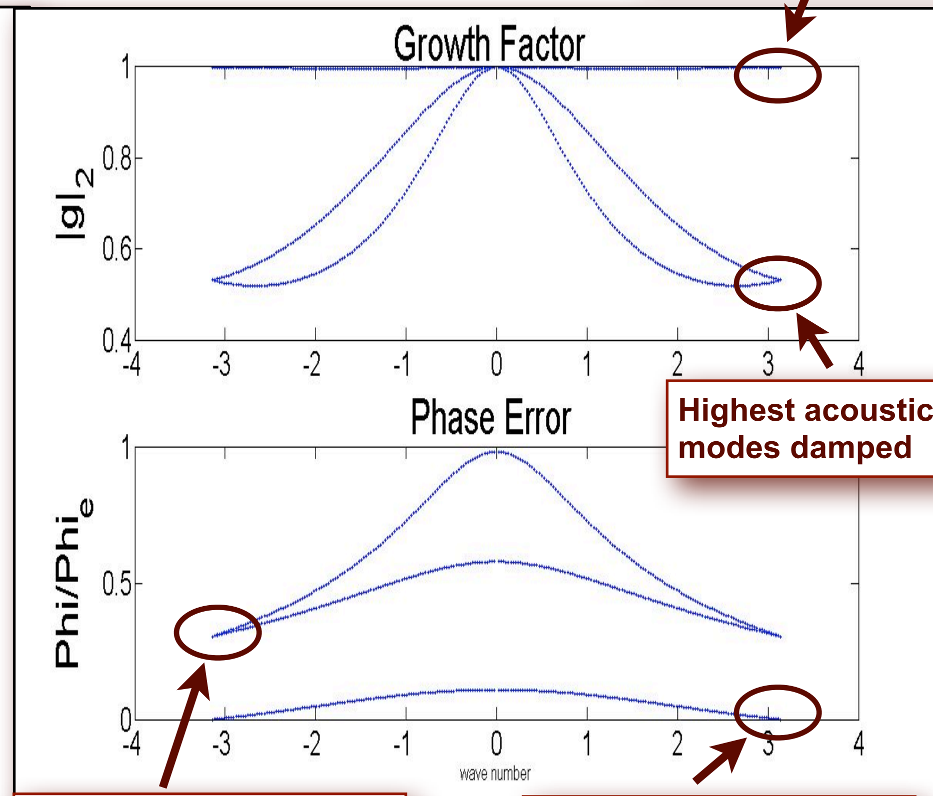
No damping of highest modes

Collocated

Staggered

Highest particle modes not damped





Highest acoustic modes damped

No convection of highest modes

Slow convection of highest acoustic modes

No convection of highest acoustic modes

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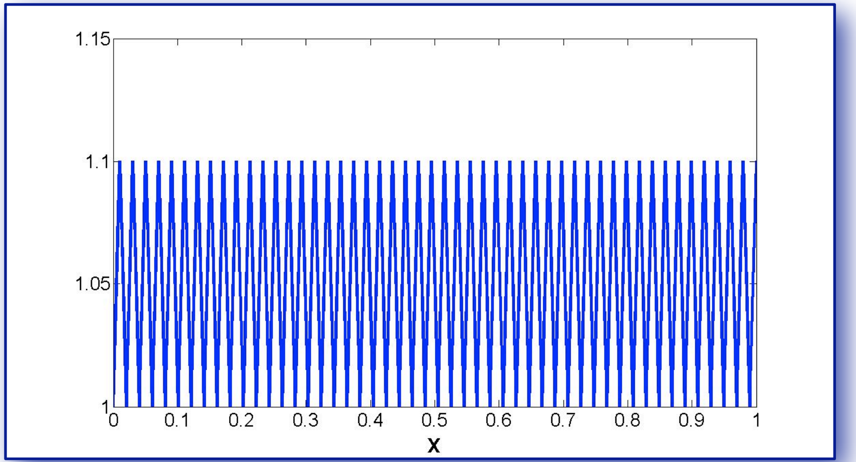
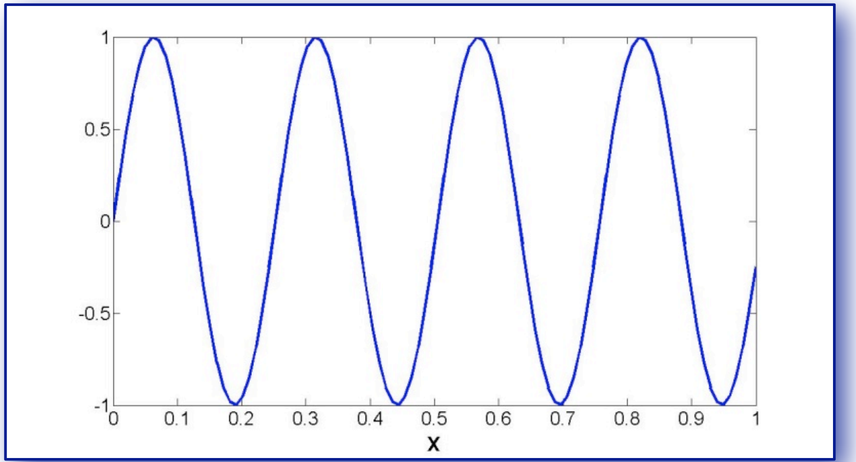
Test Cases

1D Duct
– Non-dissipative BC's $\Delta U_{IL} = \Delta U_{IL-1}$
– Periodic BC's avoid issues with reflections

Saw-tooth i.c.



Sinusoidal i.c.



For Euler Eqns:
– Use Characteristic Eqns
– However, staggered grid does not allow proper diagonalization

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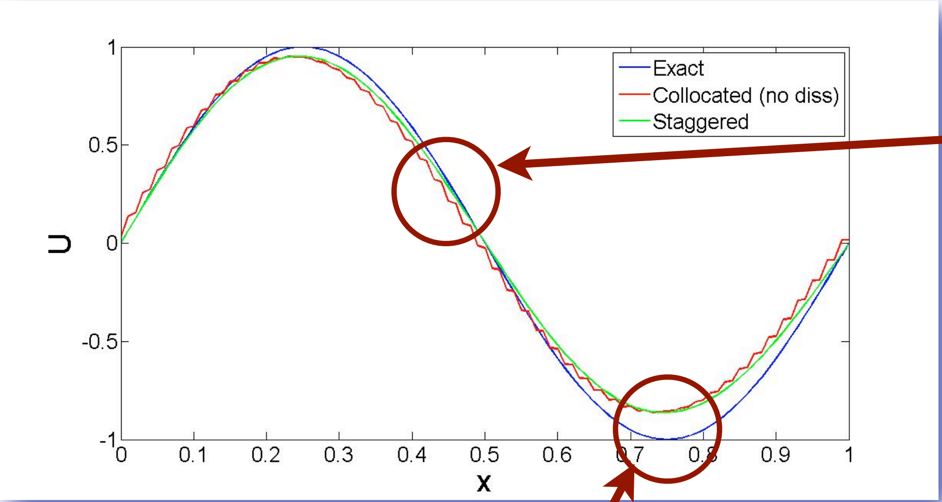
Wave Equation Results

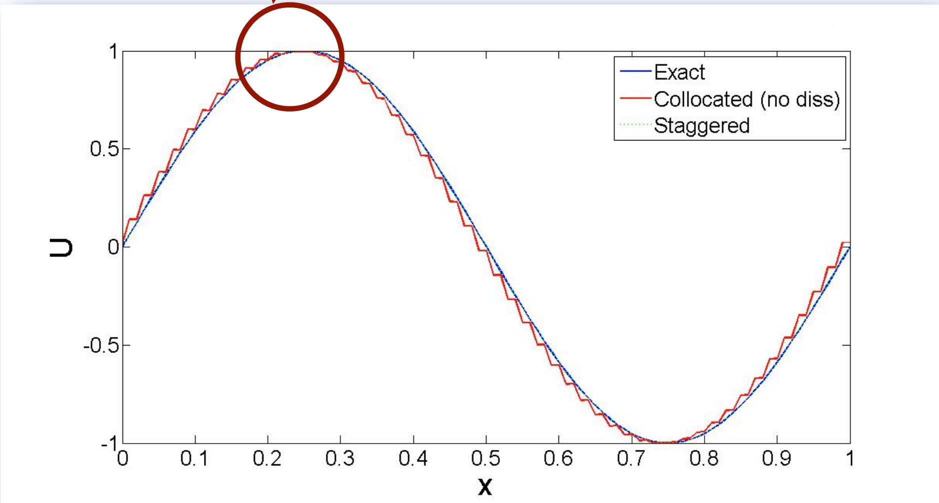
Euler Implicit



Odd-even errors in collocated grid solution; staggered solutions are smooth

Mid-wave number damping in the Euler Implicit Scheme

Crank-Nicolson



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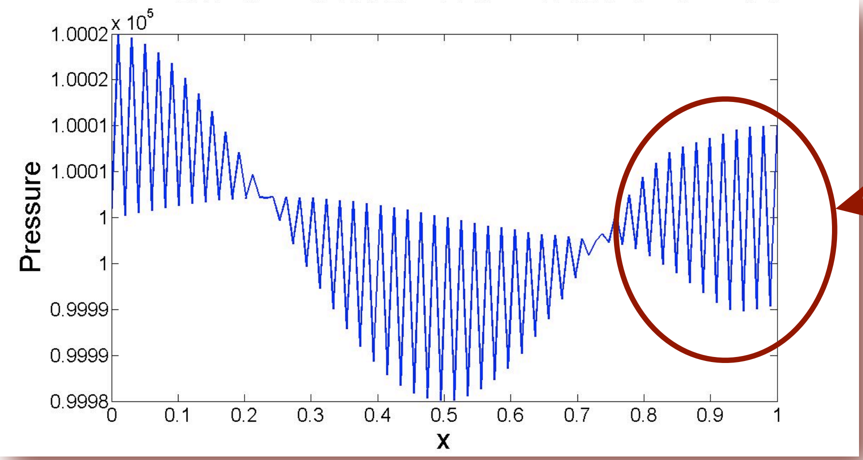
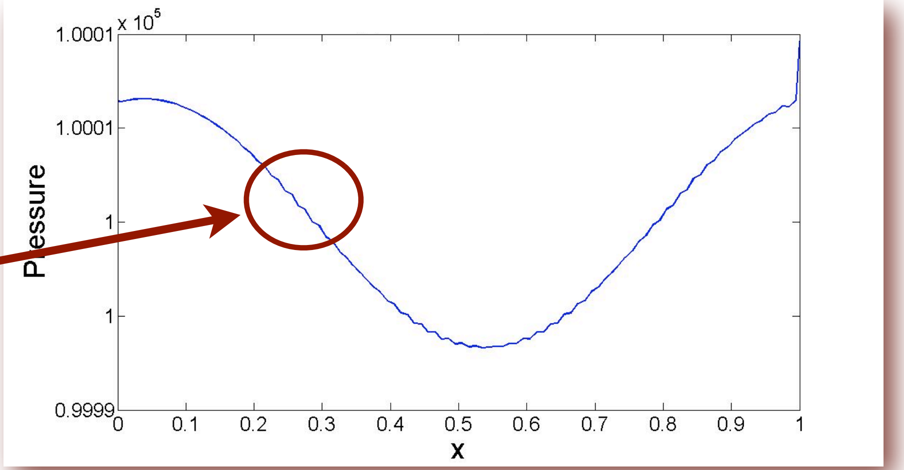
Euler Equations Results

Collocated



Collocated grid solution shows strong odd-even splitting errors

Staggered



Staggered grid solution is relatively smooth

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KE Conservative Scheme

Collocated Grid

Transport Eqn

$$\frac{\left[(\rho\phi_k)^{n+1} - (\rho\phi_k)^n\right]}{\Delta t} + \Delta_x(\rho u_j)\phi_k^* = 0$$

Time-Averaging

$$\phi_k^* = \frac{(\sqrt{\rho}\phi_k)^{n+1} + (\sqrt{\rho}\phi_k)^n}{(\sqrt{\rho})^{n+1} + (\sqrt{\rho})^n}$$

Roe-averaging in time

KE Transport Eqn

$$\frac{\left[(\rho\phi_k^2)^{n+1} - (\rho\phi_k^2)^n\right]}{2\Delta t} + \Delta_x(\rho u_j)\frac{\phi_k^2}{2} = 0$$

Ensures full KE preservation

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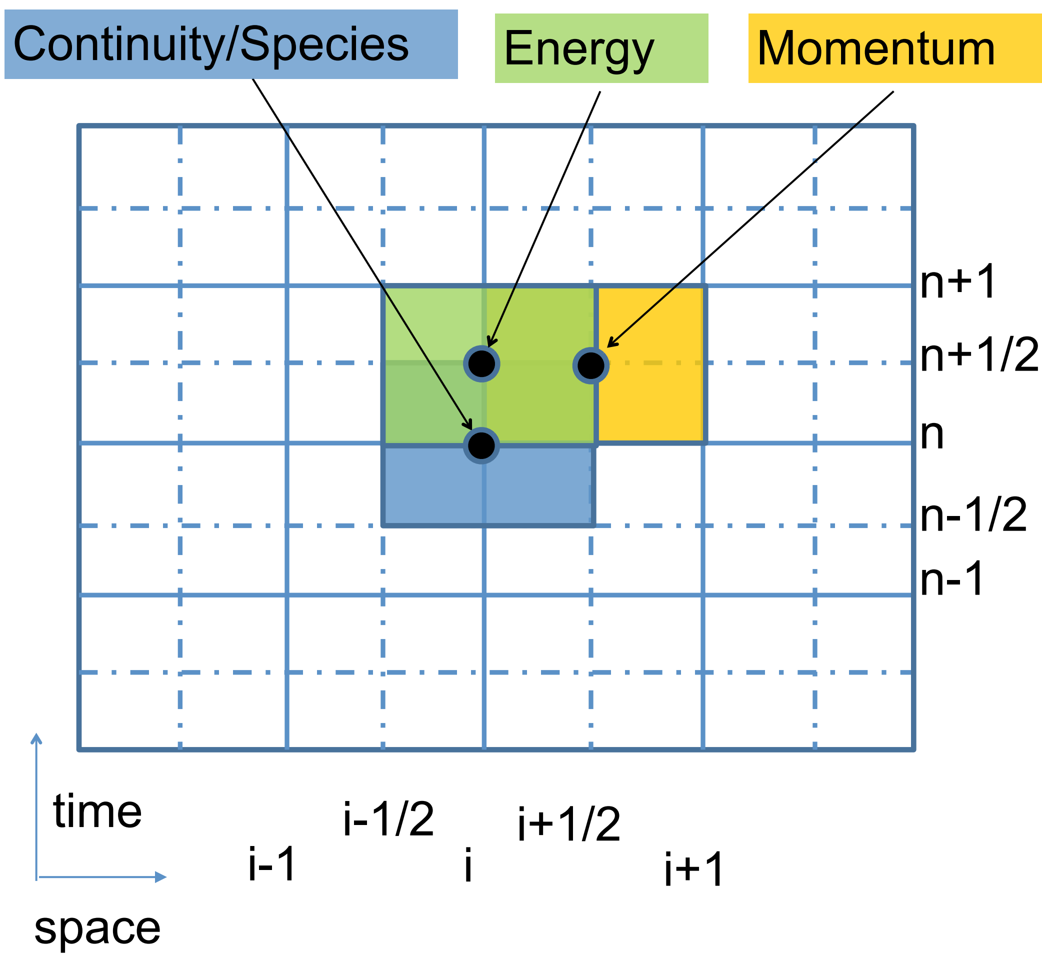
KE Conservative Scheme

Staggered Grid in Space and Time



Time-Averaging

$$(u)^*_{i+1/2,j} \equiv \frac{\left(\sqrt{\overline{\overline{\rho}^{1t}}^{1x}}\, u_\alpha\right)^{n+1}_{i+1/2,j} + \left(\sqrt{\overline{\overline{\rho}^{1t}}^{1x}}\, u_\alpha\right)^{n}_{i+1/2,j}}{\left(\sqrt{\overline{\overline{\rho}^{1t}}^{1x}}\right)^{n+1}_{i+1/2,j} + \left(\sqrt{\overline{\overline{\rho}^{1t}}^{1x}}\right)^{n}_{i+1/2,j}}$$

$$(h^0)^*_{i,j} \equiv \frac{\left(\sqrt{\overline{\rho}^{1t}}\, h^0\right)^{n+1}_{i,j} + \left(\sqrt{\overline{\rho}^{1t}}\, h^0\right)^{n}_{i,j}}{\left(\sqrt{\overline{\rho}^{1t}}\right)^{n+1}_{i,j} + \left(\sqrt{\overline{\rho}^{1t}}\right)^{n}_{i,j}}$$

$$(Y_k)^*_{i,j} = \frac{\left(\sqrt{\rho}Y_k\right)^{n+1/2}_{i,j} + \left(\sqrt{\rho}Y_k\right)^{n-1/2}_{i,j}}{\sqrt{\rho}^{\,n+1/2}_{i,j} + \sqrt{\rho}^{\,n-1/2}_{i,j}}$$

Roe-averaging in time leads to full kinetic energy preservation of momentum and scalar fields.

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Summary

- **Von Neumann Analysis provides dispersion & damping behavior**
  - Staggered grid schemes show natural damping even when artificial dissipation is **not** added explicitly
  - Dispersion errors are sometimes non-intuitive - faster wave speeds for small CFL's and slower wave-speeds for high CFL's
- **Periodic wave tests validate von Neumann results**
  - Staggered grid schemes provide smooth particle wave solutions with minimal dissipation
  - Acoustic wave damping is consequential for compressible LES
- **Kinetic Energy conservative schemes**
  - Formulated for both staggered and collocated grids
  - Schemes possess favorable properties for scalar energies
    - Maybe consequential for reacting-LES problems
  - Test results for improved schemes are forthcoming

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